

# Hydroinformatik - SoSe 2026

## UW-BHW-414-12: Grundlagen der Hydromechanik

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Dresden, 12.06.2026

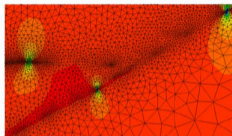
# Zeitplan: Hydroinformatik I+II

Sommersemester 2026: Stand: 06.04.2026

Nr.	KW	Datum	ID	Thema
01+02	16	17.04.2026	UW-BHW-414-01/02	Einführung in die Vorlesung, Umweltinformatik
03	16	17.04.2026	UW-BHW-414-03	Werkzeuge, Hello World (in C++)
05	17	24.04.2026	UW-BHW-414-04	Selbststudium: Software-Installationen
07	19	08.05.2026	UW-BHW-414-05	Objekt-Orientierte Programmierung: C++, Klassen
09	20	15.05.2026	UW-BHW-414-06	Programmiersprache Python
11	21	22.05.2026	UW-BHW-414-07/08	Modellierung, Digitalisierung - Wasser 4.0
00	22	29.05.2026		Vorlesungsfreie Woche
13	23	05.06.2026	UW-BHW-414-09/10	KI, Maschinelles Lernen, Neuronale Netzwerke
15	24	12.06.2026	UW-BHW-414-11/12	Kontinuumsmechanik, Hydromechanik
17	25	19.06.2026	UW-BHW-414-I	Differentialgleichungen, Näherungsverfahren
19	26	26.06.2026	UW-BHW-414-J	Finite-Differenzen, explizite Verfahren
21	27	03.07.2026	UW-BHW-414-K	Finite-Differenzen, implizite Verfahren
23	28	10.07.2026	UW-BHW-414-L	Gerinnehydraulik, Grundwasserhydraulik
25	29	17.07.2026	UW-BHW-414-M	Grundwasserhydraulik
27	30	24.07.2026	UW-BHW-414-N	Zusammenfassung, Klausurvorbereitung

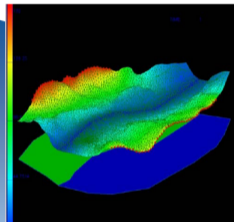
- 1 UW-BHW-414-12: Grundlagen der Hydromechanik
  - Semesterplan

$$\frac{d\psi}{dt} = \frac{\partial\psi}{\partial t} + \mathbf{v}^E \nabla\psi$$

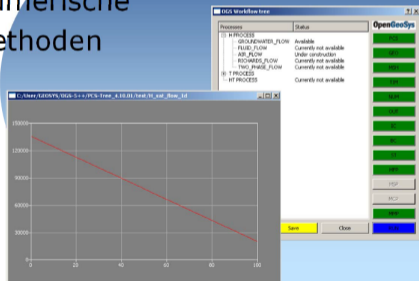


Basics  
Mechanik

Anwendung



Numerische  
Methoden



Programmierung  
Visual C++

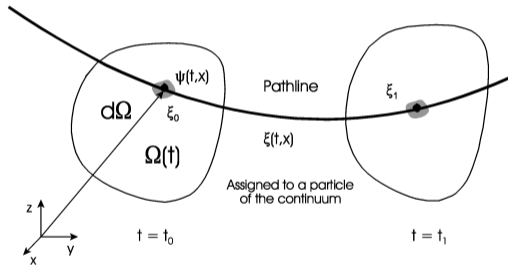
Prozessverständnis

- ▶ Erhaltungsgrößen
- ▶ Massenerhaltung
- ▶ Fluidmassenerhaltung
- ▶ Diffusion
- ▶ Impulserhaltung
- ▶ Spannungen
- ▶ Fluiddruck
- ▶ Strömungsprobleme

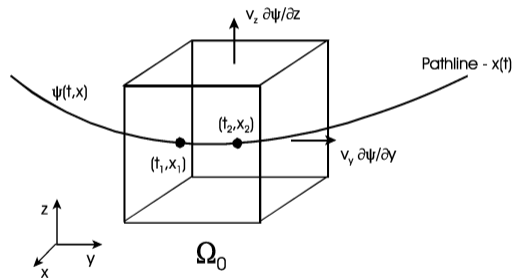
# Wiederholung

## General Balance Equation

### Lagrange



### Euler



$$\begin{aligned}\frac{d}{dt} \int_{\Omega} \psi d\Omega &= \frac{\partial}{\partial t} \int_{\Omega} \psi d\Omega + \oint_{\partial\Omega} \Phi^{\psi} \cdot d\mathbf{S} \\ &= \frac{\partial}{\partial t} \int_{\Omega} \psi d\Omega + \int_{\Omega} \nabla \cdot \Phi^{\psi} d\Omega\end{aligned}\quad (1)$$

$\lim d\Omega \rightarrow 0$

$$\begin{aligned}\frac{d\psi}{dt} &= \frac{\partial\psi}{\partial t} + \nabla \cdot \Phi^{\psi} \\ &= \frac{\partial\psi}{\partial t} + \nabla \cdot (\mathbf{v}\psi) - \nabla \cdot (\mathbf{D}^{\psi} \nabla\psi) \\ &= Q^{\psi}\end{aligned}\quad (2)$$

HA 02#2020: Komponentenschreibweise  $\nabla \cdot (\mathbf{v}\psi)$ ,  $\nabla \cdot (\mathbf{D}^{\psi} \nabla\psi)$

# Wiederholung

## Conservation Quantities (1.1.2)

The amount of a quantity in a defined volume  $\Omega$  is given by

$$\Psi = \int_{\Omega} \psi d\Omega(t) \quad (3)$$

where  $\Psi$  is an extensive conservation quantity (i.e. mass, momentum, energy) and  $\psi$  is the corresponding intensive conservation quantity such as mass density  $\rho$ , momentum density  $\rho\mathbf{v}$  or energy density  $e$ .

Extensive quantity	Symbol	Intensive quantity	Symbol
Mass	$M$	Mass density	$\rho$
Linear momentum	$\mathbf{m}$	Linear momentum density	$\rho\mathbf{v}$
Energy	$E$	Energy density	$e = \rho i + \frac{1}{2}\rho v^2$

$$\frac{d\psi}{dt} = \frac{\partial\psi}{\partial t} + \nabla \cdot (\mathbf{v}\psi) - \nabla \cdot (\mathbf{D}^\psi \nabla\psi) = Q^\psi \quad (4)$$

The differential equation of mass conservation in divergence form becomes

$$\frac{d\rho}{dt} = \frac{\partial\rho}{\partial t} + \nabla \cdot (\mathbf{v}\rho) = 0 \quad (5)$$

Partial differentiation of the above equation gives

$$\frac{d\rho}{dt} = \frac{\partial\rho}{\partial t} + \mathbf{v} \cdot \nabla\rho + \rho\nabla \cdot \mathbf{v} = 0 \quad (6)$$

## (Phase) Mass Conservation

Using the material (or convective) derivative the mass conservation equation can be rewritten as

$$\frac{\partial \rho}{\partial t} = -\rho \nabla \cdot \mathbf{v} \quad (7)$$

Note, above convective form of mass conservation equation becomes zero only for incompressible flows, i.e.

$$\frac{\partial \rho}{\partial t} = 0 \quad (8)$$

requires divergence-free flow.

$$\nabla \cdot \mathbf{v} = 0 \quad (9)$$

From eqn. (6) results that the above expression is the continuity equation for a homogeneous fluid ( $\rho = \text{const}$ ).

$$\nabla \cdot \mathbf{v} = 0 \quad (10)$$

Links:

- ▶ <https://de.wikipedia.org/wiki/Stromfunktion>
- ▶ <https://www.ingenieurkurse.de/stroemungslehre/ebene-stroemungen/quelle-und-senke-divergenz.html>

# Divergenzfreie Strömung

Übung: HyBHW-1-02-04-E1

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3 # set up a normalized grid:
4 dim= 20
5 xarray= np.arange(-dim,dim)
6 yarray= np.arange(-dim,dim)
7 # (fluid) flow from a source at L to a sink at -L:
8 L = dim/2
9 x,y = np.meshgrid(xarray,yarray)
10 vx = (x-L)/((x-L)**2+y**2) - (x+L)/((x+L)**2 +y**2)
11 vy = y/((x-L)**2+y**2) - y/((x+L)**2 +y**2)
12 # plot the flow lines:
13 plt.figure()
14 plt.quiver(x,y, vx, vy, pivot='mid')
15 plt.xlabel("$x$-axis")
16 plt.ylabel("$y$-axis")
17 plt.axis('equal')
18 plt.show()
```

**Listing:** Python code for divergence-free flow - div v

Source: University of  
Auckland

## Conservation Quantities (1.1.2)

The amount of a quantity in a defined volume  $\Omega$  is given by

$$\Psi = \int_{\Omega} \psi d\Omega(t) \quad (11)$$

where  $\Psi$  is an extensive conservation quantity (i.e. mass, momentum, energy) and  $\psi$  is the corresponding intensive conservation quantity such as mass density  $\rho$ , momentum density  $\rho\mathbf{v}$  or energy density  $e$ .

Extensive quantity	Symbol	Intensive quantity	Symbol
Mass	$M$	Mass density	$\rho$
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# Momentum Conservation

$$\psi = \rho \mathbf{v}$$

$$\frac{\partial}{\partial t} \int_{\Omega} \rho \mathbf{v} d\Omega + \oint_{\partial\Omega} \boldsymbol{\Phi}^m \cdot d\mathbf{S} = \int_{\Omega} \rho \mathbf{f} d\Omega \quad (12)$$

Flux term: The advective momentum flux is defined as

$$\boldsymbol{\Phi}^m = (\rho \mathbf{v}) \otimes \mathbf{v} = (\rho \mathbf{v}) \mathbf{v} \quad (13)$$

$$\mathbf{F} = \int_{\Omega} \rho \mathbf{f} d\Omega = \int_{\Omega} \rho (\mathbf{f}^e + \mathbf{f}^i) d\Omega = \underbrace{\int_{\Omega} \rho \mathbf{f}^e d\Omega}_{\text{External forces}} + \underbrace{\oint_{\partial\Omega} \boldsymbol{\sigma} : d\mathbf{S}}_{\text{Internal forces}} \quad (14)$$

Substituting now flux and source terms of momentum we obtain

$$\frac{\partial}{\partial t} \int_{\Omega} \rho \mathbf{v} d\Omega + \oint_{\partial\Omega} \rho \mathbf{v} (\mathbf{v} \cdot d\mathbf{S}) = \int_{\Omega} \rho \mathbf{f}^e d\Omega + \oint_{\partial\Omega} \boldsymbol{\sigma} d\mathbf{S} \quad (15)$$

Applying the Gauss-Ostrogradskian theorem to the surface integrals

$$\begin{aligned} \oint_{\partial\Omega} \rho \mathbf{v} (\mathbf{v} \cdot d\mathbf{S}) &= \int_{\Omega} \nabla \cdot (\rho \mathbf{v} \mathbf{v}) d\Omega \\ \oint_{\partial\Omega} \boldsymbol{\sigma} : d\mathbf{S} &= \int_{\Omega} \nabla \cdot \boldsymbol{\sigma} d\Omega \end{aligned} \quad (16)$$

# Momentum Conservation

The differential form of the momentum conservation law is then

$$\frac{\partial}{\partial t}(\rho \mathbf{v}) + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) = \rho \mathbf{f}^e + \nabla \cdot \boldsymbol{\sigma} \quad (17)$$

The above equation is now extended by partial integration

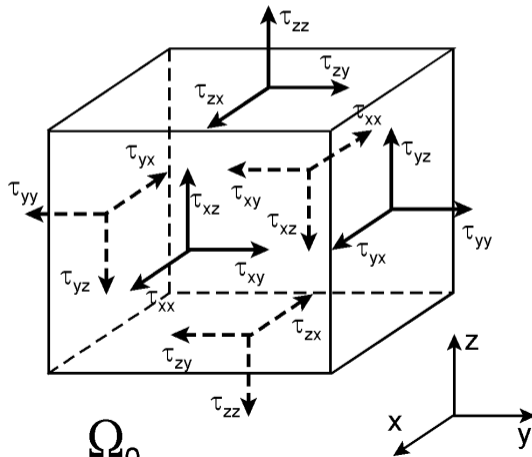
$$\begin{aligned} & \rho \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \frac{\partial \rho}{\partial t} + (\rho \mathbf{v}) \cdot \nabla \mathbf{v} + \mathbf{v} \nabla \cdot (\rho \mathbf{v}) \\ &= \rho \left[ \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right] + \mathbf{v} \left[ \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) \right] \\ & \qquad \qquad \qquad = \rho \mathbf{f}^e + \nabla \cdot \boldsymbol{\sigma} \end{aligned} \quad (18)$$

Using the mass conservation equation (5) and dividing by  $\rho$  we obtain

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}^e + \frac{1}{\rho} \nabla \cdot \boldsymbol{\sigma} \quad (19)$$

# Momentum Conservation: Stress Tensor

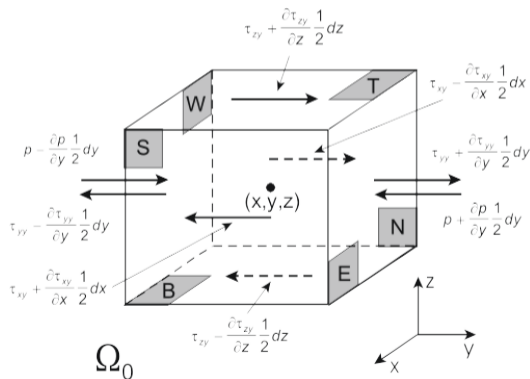
$$\sigma = -p\mathbf{I} + \tau \quad , \quad \tau = \nu \nabla \mathbf{v} \quad (20)$$



# Momentum Conservation: Stress Tensor

$$\boldsymbol{\tau} = \nu \nabla \mathbf{v}$$

(21)



$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}^e + \frac{1}{\rho} \nabla \cdot \boldsymbol{\sigma} \quad (22)$$

In index notation the above vector equation is written as

$$\begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} &= \frac{1}{\rho} \left( \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} \right) \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} &= \frac{1}{\rho} \left( \frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{yz}}{\partial z} \right) \\ \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} &= g + \frac{1}{\rho} \left( \frac{\partial \sigma_{zx}}{\partial x} + \frac{\partial \sigma_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right) \end{aligned} \quad (23)$$

with  $u = v_x, v = v_y, w = v_z$  and  $\mathbf{f}^e = \mathbf{g}$ .

# Flow Equations - Systematic

Stress Tensor

$$\boldsymbol{\sigma} = -p\mathbf{I} + \boldsymbol{\tau} \quad (24)$$

Navier-Stokes Equation

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}^e - \frac{1}{\rho} \nabla p + \nu \Delta \mathbf{v} \quad (25)$$

Euler Equation

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \mathbf{f}^e - \frac{1}{\rho} \nabla p \quad (26)$$

Stokes Equation

$$\frac{\partial \mathbf{v}}{\partial t} = \mathbf{f}^e - \frac{1}{\rho} \nabla p + \nu \Delta \mathbf{v} \quad (27)$$

Darcy Equations

$$0 = \mathbf{f}^e - \frac{1}{\rho} \nabla p + \nu \Delta \mathbf{v}$$

$$\psi = \rho_k = C_k \quad (29)$$

$$\frac{dC_k}{dt} = \frac{\partial C_k}{\partial t} + \nabla \cdot (\mathbf{v}C_k) - \nabla \cdot (\mathbf{D}_k \nabla C_k) = Q_k \quad (30)$$

# Übungen und Hausaufgaben

- EX: Übungen
- HA: Hausaufgaben

# Übungen

## Mechanik

- EX20: Jupyter Notebook
- EX21: Kontinuumsmechanik: Skalarprodukt
- EX22: Hydromechanik: Divergenzfreie Strömung
- EX23: Analytische Lösung: Elliptische Gleichung
- EX24: Analytische Lösung: Parabolische Gleichung (Diffusion)
- EX25: Analytische Lösung: Transportgleichung (ADE)
- EX26: Finite-Differenzen-Methode (FDM) explizit
- EX27: Finite-Differenzen-Methode (FDM) implizit
- EX28: Gerinnehydraulik

# Hausaufgaben

- 1 Skalarprodukt: Schreiben sie das Skalarprodukt  $\nabla \cdot \mathbf{v}$  in Komponentenschreibweise.
- 2 Mechanik: Was ist  $\mathbf{v} \cdot \nabla \psi$ ? Physikalische Bedeutung des Terms
- 3 Mechanik: Was ist  $\Phi^\psi$ ? Physikalische Bedeutung des Terms
- 4 Hydromechanik: Komponentenschreibweise  $\nabla \cdot (\mathbf{v}\psi)$
- 5 Hydromechanik: Komponentenschreibweise  $\nabla \cdot (\mathbf{D}^\psi \nabla \psi)$
- 6 Analytik: Prüfen Sie die Gültigkeit einer der Lösungen für die Diffusionsgleichung: (20), (21), (22) (siehe Vorlesung 5 >> HyBHW-S2-01-V05).
- 7 Analytik: Stellen Sie die ausgewählte analytische Lösung für die 1-D parabolische Differentialgleichung unter Verwendung der Übung EX06-parabolische-gleichung-1D.py dar. Ergänzen Sie Ihren Namen oder Matrikelnummer mit dem Befehl `plt.title("Name oder Matrikelnummer")`.
- 8 Numerik: Darstellung der numerischen Lösung (explizite FDM) für die 1-D parabolische Differentialgleichung (EX08-fdm-explicit-python). Produzieren Sie eine stabile und instabile Lösung.
- 9 Numerik: Darstellung der numerischen Lösung (implizite FDM) für die 1-D parabolische Differentialgleichung (EX09-fdm-implicit-python). Produzieren Sie die stationäre Lösung.

- zum Internet-Repository gehen (Webseite)
- Python-File editieren (Matrikel-Nummer oder Name)
- Programme zum Rechnen und Darstellen ausführen
- Ergebnis (Abbildung) in die Hausaufgaben einfügen